# Properties of MgCNi<sub>3</sub> Superconductor

#### S. MOLLAH

Department of Physics, Aligarh Muslim University, Aligarh 202 002, India E-mail: smollah@rediffmail.com

The present invited article describes the properties of superconducting MgCNi<sub>3</sub> with T<sub>c</sub> (ca. 8 K) having a simple cubic perovskite structure with space group Pm-3m and lattice parameter a ca. 3.812 Å. Its conduction electrons are derived from the partially filled Ni d-states that typically lead to the ferromagnetism in metallic Ni and many Ni-based binary alloys. It has electron type carriers in the normal state. The T<sub>c</sub> increases with the increase of x in MgC<sub>x</sub>Ni<sub>3</sub> but generally decreases due to the Ni site doping by Co, Fe, Mn, Cu etc. Again, the T<sub>c</sub> is found to increase with the increase of external pressure. The electronic contribution is slightly higher than the lattice one in its normal state thermal conductivity. The specific heat and tunneling spectroscopic studies indicate that this is an s-wave BCS-type weak/moderate coupling type-II superconductor.

Key Words: MgCNi<sub>3</sub> superconductor, Structure, Resistivity, Specific heat, Pressure effect.

#### INTRODUCTION

The MgCNi<sub>3</sub> superconductor has a perovskite structure<sup>1</sup> like CaTiO<sub>3</sub> with the equivalence of Ca to Mg, Ti to C and O to Ni. The high proportion of Ni shows that the magnetic interactions play a dominant role. It is not large enough to induce the magnetic instability in it but is associated with its superconducting properties<sup>2</sup>. <sup>13</sup>C NMR studies<sup>3</sup> suggest that the electronic states reach a modestly mass enhanced Fermiliquid like states prior to the superconducting transition. Lattice distortion associated with charge density waves (CDW) or long range antiferromagnetic (AF) ordering consistent to spin density waves (SDW) is not revealed<sup>1</sup> in MgCNi<sub>3</sub>. The single-phase perovskite structure in MgC<sub>x</sub>Ni<sub>3</sub> is found only in narrow range of carbon content (0.88 < x < 1.0). The doping of Ni site by Cu and Co decreases the T<sub>c</sub> significantly<sup>4,5</sup>. The present article discusses the synthesis and properties of MgCNi<sub>3</sub> superconductor.

#### EXPERIMENTAL

## Synthesis of MgCNi<sub>3</sub>

The MgC<sub>x</sub>Ni<sub>3</sub> samples with x = 0.9-1.5 are prepared by using the raw materials as Mg flakes, fine Ni powder and glassy carbon spherical powder<sup>1,6,7</sup>. The starting materials are properly mixed and pressed into pellets. The pellets are placed into a Ta foil, put in an alumina boat and fired in a quartz tube furnace under a mixed gas of 95% Ar and 5% H<sub>2</sub> environment. The samples are heat treated at 600°C for 0.5 h and followed by one hour at 900°C. These are cooled, grounded,

3266 Mollah Asian J. Chem.

pressed and heated at 900°C for one more hour. Owing to the volatility of Mg, 20% excess of its stoichiometric ratio is added to the initial mixture<sup>1,7</sup>. Chemical doping as well as diffusion of metallic particles in MgCNi<sub>3</sub> has been reported by several groups<sup>7</sup>. Most of the groups dope on the Ni site<sup>7</sup>. However, the doping effect on Mg site is also studied<sup>7</sup>. The preparation technique is over all conventional type.

# RESULTS AND DISUCSSION

Structure of MgCNi<sub>3</sub>

MgCNi<sub>3</sub> superconductor possesses the classical cubic perovskite structure with the space group Pm-3m and the lattice constant<sup>1,2</sup> a ca. 3.81221 Å at 295 K. The neutron diffraction study<sup>1,2</sup> shows that the formula for the superconducting phase is MgC<sub>0.96</sub>Ni<sub>3</sub> for the nominal composition MgC<sub>1.25</sub>Ni<sub>3</sub>. This is due to the small amount of unreacted graphite found in the sample<sup>1,2</sup>. The atomic positions are: Mg 1a (0, 0, 0); C 1b (0.5, 0.5, 0.5) and Ni 3c (0, 0.5, 0.5) respectively with temperature factors 0.90(3), 0.54(4) and 0.75(1) Å<sup>2</sup>. Regular domains, with an average size as small as ca. 4 nm, appear commonly in the superconducting phase<sup>7</sup>. Introduction of carbon vacancies has significant effects on the positions of the Ni atoms. Electronic structure of MgC<sub>1-x</sub>Ni<sub>3</sub> with X-ray photoemission spectroscopy and X-ray absorption spectroscopy show that overall band structure is in reasonable agreement with band structure calculations<sup>8</sup> including the existence of von Hove singularity near  $E_F$ .

In MgCNi<sub>3-x</sub>Co<sub>x</sub> system<sup>5</sup>, the lattice parameter decreases slightly with increasing x. The *in situ* high-pressure energy dispersive X-ray diffraction has shown that the structure of MgCNi<sub>3</sub> is stable<sup>7</sup> under a pressure ca. 22 GPa. MgCNi<sub>3</sub> has the perovskite structure over the whole temperature range and no structural or long-range magnetic ordering transitions are observed. There are no unusual changes of the structural

parameters near T<sub>c</sub>.

### Properties of MgCNi<sub>3</sub>

### 1) Electrical properties

The upper part of Fig. 1 shows the temperature variation of resistivity for  $MgC_xNi_3$  (x = 1-1.5) samples. Superconducting transition temperature ( $T_c$ ) is found to be ca. 6.5-7.7 K depending on the values of x. Thin films thicker than ca. 40 nm have a  $T_c$  ca. 8 K which is comparable to that of polycrystalline bulk samples. However, the  $T_c$  decreases with the decrease of film thickness<sup>7</sup>.

 $T_c$  of MgC<sub>x</sub>Ni<sub>3</sub> (x = 1.0-1.5) increases with the increase of x (Fig. 1). The highest  $T_c$  corresponds to x ca. 1.45-1.5 and decreases with further increase of x. The  $T_c$  is found to decrease systematically with decreasing

carbon concentration from the stoichiometric value<sup>1</sup>. Excess of Mg and C in initial material mixture is favourable to improve the T<sub>c</sub> and to obtain the single-phase samples<sup>1,5</sup>. Doping at Ni site by Co, Fe, Mn, Cu etc. also shows the decrease of T<sub>c</sub> except an initial increase with Fe doping<sup>9</sup>. In addition, Das and Kremer<sup>10</sup> observe a rapid suppression of superconductivity (ca. -21 K/at% Mn) in Mn substituted MgCNi<sub>3</sub>.

Yang et al.<sup>6</sup> observe the increase of  $T_c$  with pressure (P) from ac susceptibility measurement of three  $MgC_xNi_3$  (Fig. 2) with  $dT_c/dP$  ca. 0.015 K/kbar. The change of  $T_c$  with the unit cell volume (V) can be given by  $V_c/T_c/dT_c/dV = dlnT_c/dlnV = -(B/T_c)(dT_c/dP)$ , where B is the bulk modulus of the superconductor. The  $dlnT_c/dlnV$  values vary from -3.18 to -2.58 similar to  $MgB_2$  superconductor (+ 4.16) with opposite sign. The  $T_c$  can be expressed by the McMillan formula  $T_c = (\theta_D/1.45) \exp\{-1.04(1+\lambda)/[\lambda-\mu^*(1+0.62\lambda)]\}$ , where,  $\mu^*$  is the Coulomb pseudo potential,  $\theta_D$  is the Debye temperature and  $\lambda$  is the electron-phonon coupling constant. The change of  $\lambda$  and  $\theta_D$  by pressure determines of the sign of  $dT_c/dP$ . The positive  $dT_c/dP$  for  $MgC_xNi_3$  is possibly originated from the increase of  $N(E_F)$  and consequently by the enhancement of electron-phonon coupling constant  $\lambda$  if  $\mu^*$  and  $ext{-1}$  are less pressure dependent.

# 2) Magnetic properties

T<sub>c</sub> of MgCNi<sub>3</sub> decreases with the increase of magnetic field and finally becomes non-superconducting<sup>7</sup>. A magnetic field of 8 T leads to about 50% suppression of T<sub>c</sub> and a complete suppression takes place at 14 T. The width  $\Delta H = H_{90} - H_{10}$ , with  $H_{90}$  and  $H_{10}$  being respectively, the field values where 90 and 10% of the normal state resistivity is observed, remains constant ca. 0.6 T down to low temperature. This indicates that MgCNi<sub>3</sub> has a small anisotropy in H<sub>c2</sub> as the strongly anisotropic superconductor shows a gradual broadening of the superconducting transition with the decrease of temperature. The relation  $H_{c2}(0) \approx 0.69 T_c (dH_{c2}/dT)_{Tc}$  leads to the  $H_{c2}(0)$  ca. 13.2  $\pm$  0.7 T. It is found that the physical properties of MgCNi3 are very similar to those of Nb<sub>0.5</sub>Ti<sub>0.5</sub>. Thus both the compounds may have a similar relation between H<sub>c2</sub>(0) and (dH<sub>c2</sub>/dT)<sub>Tc</sub>. It is contended that MgCNi<sub>3</sub> has a Werthamer-Helfand-Hohenberg (WHH) like temperature dependence of  $H_{c2}(T)$  and the quadratic relationship  $H_{c2}(0) = 0.0237(1+\lambda)^{2.2}T_c^2/(10^5 \text{ x})$ v<sub>F</sub><sup>2</sup>), with v<sub>F</sub> as the bare Fermi velocity which points to an effective predominant single band behaviour near the quasi clean limit. The Paulilimiting field  $H_p(0) = 1.84 \times 10^4 T_c$  is expected within the weak-coupling BCS theory. The obtained  $H_{c2}(0)$  from WHH theory is higher than  $H_{P}(0)$ , suggesting that pair-breaking effects due to the Zeeman energy in MgCNi<sub>3</sub> is small<sup>7</sup>. For type-II superconductors, H<sub>p</sub>(0) should satisfy the 3268 Mollah Asian J. Chem.

relation  $H_{c2}(0) \le H_p(0)$ . The  $H_p(0)$  of MgCNi<sub>3</sub> is of the order of its  $H_{c2}(0)$  indicating the type-II superconductivity. Thermodynamic critical field  $H_c(0)$  ca. 0.18-0.6 T, Ginzburg-London (GL) coherence length  $\xi_{GL}(0)$  ca. 45-56 Å, penetration depth  $\lambda_{GL}(0)$  ca. 1800-2480 Å, and lower critical field  $H_{c1}(0)$  ca. 10-12.6 mT and  $\kappa(0)$  [=  $\lambda_{GL}(0)/\xi_{GL}(0)$ ] ca. 43.3-66 are found for this superconductor.<sup>7</sup> All the obtained parameters  $\xi_{GL}(0)$ ,  $\lambda_{GL}(0)$ , and  $\kappa(0)$  of MgCNi<sub>3</sub> also satisfy the conditions for type-II superconductivity.

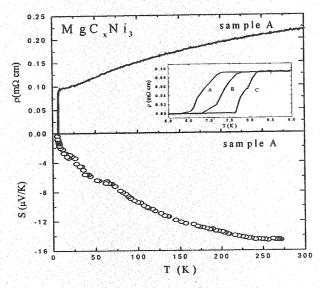


Fig. 1. Temperature (T) variation of resistivity ( $\rho$ ) and ac magnetic susceptibility ( $\chi_{ac}$ ) thermoelectric power (S) for sample A (MgCNi<sub>3</sub>) at samples6, 7 with x = 1.0 (A), 1.25 (B) ambient pressure. The inset shows the resistivity ( $\rho$ ) of at various pressures (P). the three MgC<sub>x</sub>Ni<sub>3</sub> samples<sup>6,7</sup> with x = 1.0 (A), x = 1.25 (B) and x= 1.5 (C) near T<sub>c</sub>

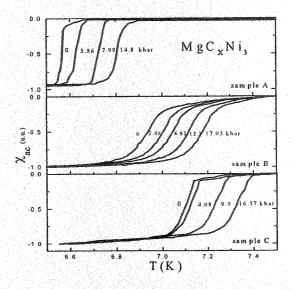


Fig. 2. Variation of of MgC<sub>x</sub>Ni<sub>3</sub> and 1.5 (C) near T<sub>c</sub>

The critical current density (J<sub>c</sub>) of MgC<sub>1.5</sub>Ni<sub>3</sub> ~  $10^3$ - $10^4$  A/cm<sup>2</sup> at 4.2 K. Normal state magnetoresistance (MR) is given by  $\Delta \rho/\rho_0 = [\rho(H)$ -

 $ho_0]/
ho_0$  where ho(H) and  $ho_0$  are respectively the resistivities with or without a magnetic field H. The resistance of MgCNi<sub>3</sub> increases with magnetic filed showing a positive magnetoresistance. However, the MR of MgCNi<sub>3</sub> is much smaller than the borocarbide superconductors. The Hall coefficient (R<sub>H</sub>) of MgCNi<sub>3</sub> at a magnetic filed of 10 T is almost constant up to a temperature of 140 K and beyond that the magnitude of it decreases with the increase of temperature. R<sub>H</sub> of MgCNi<sub>3</sub> is negative for the whole temperature range that definitely indicates that the carrier in MgCNi<sub>3</sub> is electron type and is supported by its thermoelectric power data. At T = 100 K, R<sub>H</sub> =  $-6.1 \times 10^{-10}$  m<sup>3</sup>/C and the carrier density (n) is  $ca. 1.0 \times 10^{22}$ /cm<sup>3</sup>, which is comparable with that of the theoretically calculated value ( $1.3 \times 10^{22}$ /cm<sup>3</sup>) and that in perovskite (Ba,K)BiO<sub>3</sub>, but less than that of the metallic binary MgB<sub>2</sub>.

### 3) Thermal properties

Thermal conductivity (k) of MgCNi<sub>3</sub> is nearly constant<sup>7</sup> above 210 K. It is of the order of intermetallics, larger than that of borocarbides and smaller than MgB<sub>2</sub>. It has been observed that the electronic contribution is slightly higher than the lattice contribution in the normal state. It has been also found from thermal conductivity that the scattering of electrons with static imperfections of the crystal becomes dominant near  $T_c$ . Lower part of Fig. 1 shows the thermoelectric power (S) of MgC<sub>x</sub>Ni<sub>3</sub> sample with x = 1.0 (sample A). The temperature dependence of S is negative confirming the carriers to be electron type, which is consistent with other published results and inconsistent with the theoretical predictions<sup>7</sup>. The absolute value of S(275 K) ca. 9.5-14  $\mu$ V/K of this sample decreases with decreasing temperature and is a characteristic of metallic transport in normal state<sup>6,7</sup>. The magnetic field dependence of specific heat (C) suggests that MgCNi<sub>3</sub> is an s-wave superconductor in nature<sup>7</sup>.

### 4) Energy gap

The superconducting energy gap  $\Delta$  is found<sup>7</sup> to be ca. 1.1-1.15 meV from different measurements. These values of  $\Delta$  are consistent with the calculated one from the specific heat and other data. The tunneling spectroscopy as well as the specific heat studies show that  $2\Delta/k_BT_c$  of MgCNi<sub>3</sub> varies from 3.75 to 5 which is higher than the typical weak coupling BCS value (ca. 3.52). The electron-phonon coupling constant  $\lambda$  = 0.66-0.84 indicates the moderate coupling superconductivity in MgCNi<sub>3</sub>. Therefore, it can be concluded that MgCNi<sub>3</sub> is a BCS s-wave, moderate coupling, type-II and single gap superconductor.

#### Conclusions

The synthesis, structure and properties of MgCNi<sub>3</sub> superconductor are discussed. It has a perovskite structure. T<sub>c</sub> of MgC<sub>x</sub>Ni<sub>3</sub> is sensitive to carbon content increasing with x. However, the doping on Mg and Ni site decreases the T<sub>c</sub>. External pressure increases the T<sub>c</sub> of MgC<sub>x</sub>Ni<sub>3</sub>. Hall coefficient and thermoelectric power data of MgCNi<sub>3</sub> show that the carriers in this compound are electrons. The density of states (DOS) of the Fermi level (E<sub>F</sub>) is dominated by Ni 3d-states and there is a von Hove singularity of the DOS just below the E<sub>F</sub>. This is a single gap superconductor in contrast to MgB<sub>2</sub>. The field dependent specific heat and resistivity results imply that it is a moderate coupling, type-II and swave BCS superconductor.

## REFERENCES

- 1. T. He, Q. Huang, A.P. Ramirez, Y. Wang, K.A. Regan, N. Rogado, M.A. Hayward, M.K. Hass, J.S. Slusky, K. Inumaru, H.W. Zandbergen, N.P. Ong and R.J. Cava, *Nature (London,)* 411, 54 (2001).
- 2. Q. Huang, T. He, K.A. Regan, N. Rogado, M. Hayward, M.K. Hass, K. Inumaru and R.J. Cava, *Physica C*, 363, 215 (2001).
- 3. P.M. Singer, T. Imai, T. He, M.A. Hayward and R.J. Cava, *Phys. Rev. Lett.*, 87, 257601 (2001).
- 4. M.A. Hayward, M.K. Hass, A.P. Ramirez, T. He, K.A. Regan, N. Rogado, K. Inumaru and R.J. Cava, *Solid State Commun.* 119, 491 (2001).
- 5. Z.A. Ren, G.C. Che, S.L. Jia, H. Chen, Y.M. Ni and Z.X. Zhao, Sci. Chin. Ser. A-Math. Phys. Astron., 44, 1205 (2001).
- H.D. Yang, S. Mollah, W.L. Huang, P.L. Ho, H.L. Huang, C.-J. Liu, J.-Y. Lin, Y.-L. Zhang, R.-C. Yu and C.-Q. Jin, *Phys. Rev.* B, 68, 092507 (2003).
- 7. S. Mollah, J. Phys.: Condens. Matter, 16, R1237 (2004).
- 8. J.H. Kim, J.S. Ahn, J. Kim, M.-S. Park, S.I. Lee, E.J. Choi and S.-J. Oh, *Phys. Rev.* B, 66, 172507 (2002).
- 9. T.G. Kumary, J. Janaki, A. Mani, S.M. Jaya, V.S. Sastry, Y. Hariharan, T.S. Radhakrishnan and M.C. Valsakumar, *Phys. Rev.* B **66**, 064510 (2002).
- 10. Das and R.K. Kremer, Phys. Rev. B, 68, 064503 (2003).