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# Magnetic Properties of Ln<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub> (Ln = Ho and Er) Compound

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The heavy rare earth zirconate compound  $Ln_2Zr_2O_7$  (where Ln = Ho, Er) were prepared by solid state reaction method. The prepared compounds have been characterized by XRD, DT, TG and DTG. XRD studies suggest the formation of single phase orthorhombic structure at room temperature. We describe magnetic susceptibility measurements of the compounds at high temperature. In these compounds magnetism arises from the rare-earth ions, while Zr is non magnetic. All the compounds show antiferromagnetic behaviour at low temperature and obey Curie-Weiss law.

Key Words: Magnetic susceptibility, XRD, Thermal properties, Ho<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub>, Er<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub>.

### **INTRODUCTION**

The compounds with general formula  $A_2B_2O_7$ , where A is trivalent rare earth and B is tetravalent transition element, have attracted a lot of attention in recent years<sup>1-3</sup>. However, very few studies have been performed on 4d transition metal oxides. This seems to be due to the peculiar magnetic properties of the 4d transition metal ions. In general, since 4d electrons have a large spatial extent, very strong spin-orbit coupling and a large ligand-field effect are expected<sup>4</sup>. When the energy splitting caused by spin-orbit coupling is equivalent to that of the multiplets, a higher energy term offen mixes into the lowest lying term. For this reason, 4d transition metal ions offen have extremely small magnetic susceptibilities that caused be explained by hunds rule, which explained the magnetism of iron-group ions. This seems to make the study on oxides with 4d transition metal ions difficult. However, since 4d transition metal oxides are very interesting candidates for use as metallic conductors, catalysts and materials of photoelectrolytic cells and superconductors. It seems important to investigate their magnetic properties in details.

Several such compounds have been investigated in the past<sup>5-11</sup>. In this paper, we present the preparation, characterization and magnetic susceptibilities of  $Ho_2Zr_2O_7$  and  $Er_2Zr_2O_7$ .

#### **EXPERIMENTAL**

The starting material for the preparation of  $Ln_2Zr_2O_7$ (where Ln = Ho and Er) compounds were Ho<sub>2</sub>O<sub>3</sub> ( $\approx$  99.9 % purity, Alfa Aesar), Er<sub>2</sub>O<sub>3</sub> ( $\approx$  99.99 % purity, Alfa Aesar) and ZrO<sub>2</sub> ( $\approx$  99.7 % purity, Alfa Aesar). The stoichiometric amount of these oxides were mixed and heated in an alumina crucible for 50 h at a temperature of 1300 K followed by one intermediate grinding and final product was cool down slowly. The cooled material was powdered and the formation of the prepared material was checked by X-ray diffraction technique at room temperature.

The X-ray diffraction of the compounds was studied using X-ray diffractometer (Thermoelectron-ARL EXTRA) at room temperature by using CuK<sub> $\alpha$ </sub> radiation with  $\lambda = 0.15418$  nm in a wide range of Bragg angle ( $10^{\circ} \le 2\theta \le 80^{\circ}$ ). The DT, TG and DTG studies of the compounds were carried out in nitrogen gas using a thermal analyzer (Perkin-Elmer Pyris) at a heating rate of 283 K/min and flow rate of 100 mL/min from 323 to 1123 K. Magnetic susceptibility measurements were done on powdered samples using Faraday's method<sup>12,13</sup>. Gd<sub>2</sub>(WO<sub>4</sub>)<sub>3</sub> has been used for standardization.

## **RESULTS AND DISCUSSION**

The X-ray diffraction pattern of the compound was taken at room temperature using  $CuK_{\alpha}$  line ( $\lambda = 0.15418$  nm) as shown in Figs. 1 and 2. From XRD pattern,  $d_{hkl}$  values have been evaluated using relation<sup>14</sup>.

$$d_{hkl} = \frac{0.15418}{2\text{Sin}\theta} \tag{1}$$

From these values of  $d_{hkl}$ , structures of the studied compounds were resolved using usual procedure. All the peaks have been assigned with proper hkl planes. This confirms that prepared compounds has single phase and no unreacted part of the starting material was left. The unit cell parameters are given in Table-1.









Fig. 2. Room temperature XRD pattern of Er<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub>

TABLE-1 STRUCTURAL PARAMETERS OF Ln <sub>2</sub> Zr <sub>2</sub> O <sub>7</sub> HAVING ORTHORHOMBIC UNIT CELL			
Compound	Lattice parameters		
	$a_0 (nm)$	b <sub>0</sub> (nm)	$c_0 (nm)$
$Ho_2Zr_2O_7$	0.8676	0.7394	0.6132
$Er_2Zr_2O_7$	1.2903	0.9123	0.3633

The differential thermal (DT), thermogravimetric (TG) and derivative thermogravimetry (DTG) thermogram of  $Ln_2Zr_2O_7$  compounds are shown in Figs. 3 and 4. The DT thermogram of  $Er_2Zr_2O_7$  show exothermic and endothermic peaks at 353 and 373 K, respectively.



Fig. 3. DT, TG and DTG thermogram of Ho<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub>

In TG thermogram, the compound losses mass in successive steps. The mass loss and temperature region is shown in Table-2 for  $Ho_2Zr_2O_7$  and  $Er_2Zr_2O_7$ .



Fig. 4. DT, TG and DTG thermogram of Er<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub>

TABLE-2 MASS LOSS (%) CALCULATED FROM TG THERMOGRAM				
	Mass loss (%) in the			
Compound	1 <sup>st</sup> step		2 <sup>nd</sup> step	
compound	Temp.	Mass loss	Temp.	Mass loss
	region (K)	(%)	region (K)	(%)
$Ho_2Zr_2O_7$	323-548	0.02	-	-
$Er_2Zr_2O_7$	323-498	0.05	498-823	0.37

The first step of mass loss may be due to removal of absorbed water and other gaseous species and the second step of mass loss could be responsible for major thermal dehydration of the compound. DTG thermogram of  $Ho_2Zr_2O_7$  and  $Er_2Zr_2O_7$  show maximum rate of mass change at 528 and 635 K, respectively.

The magnetic susceptibility measurement of all the studied compound was done in heating and cooling cycle. No hysteresis was observed and  $\chi_M$  values were found to be same in heating and cooling cycles. Although a small loss of weight is detected for all the compounds in heating cycle may be due to presence of moisture. The variation  $\chi_M^{-1}$  vs. T of compounds in the temperature range 300-1100 K at field of  $1.55 \times 10^{-1}$  Tesla are shown in Figs. 5 and 6. It is seen from these plots that at higher temperature the  $\chi_M^{-1}$  vs. T plots are linear and obey Curie-Weiss law and expressed by the relation<sup>15</sup>.

$$\chi_{\rm M}^{-1} = \frac{T - \theta_{\rm p}}{\overline{\rm C}_{\rm M}} \tag{2}$$

where  $\theta_p$  is paramagnetic Curie temperature and  $\overline{C}_M$  is the molar Curie constant.

All the studied compounds are magnetically simple because magnetism arises from the trivalent rare earth ions *i.e.* the magnetic interaction exists in these compounds is  $Ln^{3+}$ -  $Ln^{3+}$ . Thus at temperature much higher than ordering temperature the molar magnetic susceptibility of all these compounds can be approximated by the relation<sup>16</sup>.

$$\chi_{\rm M} = \frac{N\mu_0\mu_\beta^2}{3k} \left[ \frac{\overline{p}^2}{T - \theta_p} \right]$$
(3)

where N is Avagadro number,  $\mu_B$  is Bohr magneton,  $\mu_0$  is permeability constant, k is Boltzmann constant,  $\overline{p}$  magneton numbers of magnetic ions  $Ln^{3+}$  and  $\theta_p$  is the paramagnetic Curie temperature.



Fig. 5. Variation of inverse of molar magnetic susceptibility  $(\chi_M^{-1})$  with absolute temperature of  $Ho_2Zr_2O_7$ 



Fig. 6. Variation of inverse of molar magnetic susceptibility  $(\chi_M^{-1})$  with absolute temperature of  $Er_2Zr_2O_7$ 

We can write above equation as:

$$\chi_{\rm M}^{-1} = \frac{3k(T - \theta_{\rm p})}{N\mu_0\mu_{\rm B}^2\overline{\rm p}^2} \tag{4}$$

Comparing eqns. 2 and 4 we have

$$\overline{C}_{M} = \frac{N\mu_{0}\mu_{\beta}^{2}\overline{p}^{2}}{3k}$$

This yield

$$\overline{p} = \left[\frac{3k\overline{C}_{M}}{N\mu_{0}\mu_{\beta}^{2}}\right]^{1/2}$$
(5)

The experimental value of  $\overline{p}$  can be evaluated from the value of  $\overline{C}_M$  obtained from  $\chi_M^{-1} vs$ . T plot. The theoretical value of  $\overline{p}$  has been already known. The theoretical and experimental values of  $\overline{p}$  with magnetic ions are given in Table-3 and the values of  $\theta_p$  and  $\overline{C}_M$  are given in Table-4.

It is seen from the table that there is a good agreement between theoretical and experimental values of  $\overline{p}$ , which shows that ionic moment involved in the magnetization process concern the tripositive rare earth ions.

TABLE-3
MAGNETIC ION WITH THEORETICAL AND EXPERIMENTAL
VALUE OF AVERAGE MAGNETON NUMBER ( $\overline{p}$ ) OF THE
In 7r O

Compounds	Magnetic ion	Theoretical value	Experimental value
$Ho_2Zr_2O_7$	Ho <sup>3+</sup>	10.61	10.46
$Er_2Zr_2O_7$	Er <sup>3+</sup>	9.58	9.50

$\begin{array}{c} \mbox{PARAMAGNETIC CURIE TEMPERATURE } (\theta_{\rm p}) \mbox{ AND MOLAR} \\ \mbox{ CURIE CONSTANT } (\box{$\overline{C}_M$}) \mbox{ OF } \mbox{Ln}_2 \mbox{Zr}_2 \mbox{O}_7 \\ \hline \box{ Compounds } \box{$\theta_{\rm p}$} (K) \box{$\overline{C}_M$} \times 10^5 \mbox{ (m}^3 \mbox{ K mol}^{-1}) \\ \hline  \end{array}$	TABLE-4			
$\begin{array}{c} \text{CURIE CONSTANT}  ( \overline{C}_{M} )  \text{OF}  \text{Ln}_{2} \text{Zr}_{2} \text{O}_{7} \\ \hline \\ \hline \text{Compounds} \qquad \theta_{_{P}}  (\text{K}) \qquad \overline{C}_{M} \times 10^{5}  (\text{m}^{3}  \text{K}  \text{mol}^{-1}) \end{array}$	PARAMAGNETIC CURIE TEMPERATURE $(\theta_p)$ AND MOLAR			
Compounds $\theta_{p}(K) = \overline{C}_{M} \times 10^{5} (m^{3} K \text{ mol}^{-1})$	CURIE CONSTANT ( $\overline{C}_M$ ) OF Ln <sub>2</sub> Zr <sub>2</sub> O <sub>7</sub>			
	Compounds	$\theta_{p}(K)$	$\overline{C}_{M} \times 10^{5} (\text{m}^{3}\text{K mol}^{-1})$	
Ho <sub>2</sub> Zr <sub>2</sub> O <sub>7</sub> -25 17.11	Ho <sub>2</sub> Zr <sub>2</sub> O <sub>7</sub>	-25	17.11	
$Er_2Zr_2O_7$ -10 14.11	$Er_2Zr_2O_7$	-10	14.11	

The values of  $\theta_p$  are negative for studied compounds suggesting a possible antiferromagnetic ordering of these compounds at lower temperature. However, such small values of  $\theta_p$  can also be due to crystal field effect with a little contribution from simple dipole-dipole interaction between the magnetic ions.

#### Conclusion

From XRD studies, it is confirmed that the studied compounds have single phase orthorhombic structure at room temperature. DT, TG and DTG studies show that the compounds are stable above certain temperature. The high temperature magnetic measurement shows the magnetic ions contribute towards magnetic susceptibility as per their effective magneton number. All the studied compounds obey Curie-Weiss law behaviour at higher temperature. They have negative value of paramagnetic Curie temperature indicating antiferromagnetic ordering at lower temperature.

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#### REFERENCES

- 1. K.J. Standley, Oxide Magnetic Materials, Oxford, Clarendon Press (1972).
- 2. E.P. Wohlforth, Hands Book of Magnetically Ordered Materials, Amsterdam, North Holland I & II (1980).
- E.P. Wohlforth, A Hand Book on the Properties of Magnetically Ordered Substances, Ferromagnetic Materials, Amsterdam, North Holland III (1982).
- 4. J.H. Van Vleck, The Theory of Electric and Magnetic Susceptibilities, Oxford University Press, Oxford, p. 311 (1965).
- 5. S.T. Bramwell, M.N. Field, M.J. Harris and J.P. Parkin, *J. Phys. Condens. Matter*, **12**, 483 (2000).
- Q. Xu, W. Pan, J. Wang, C. Wan, L. Qi and H. Miao, J. Am. Ceram. Soc., 89, 340 (2006).
- Y.X. Qin, J.D. Wang, W. Pan, C.L. Wan and Z.X. Qu, *Key Eng. Mater.*, 336, 1764 (2007).
- Y.P. Tong, Z.X. Yu, L.D. Lu, X.J. Yang and X. Wang, *Mater. Res. Bull.*, 43, 2736 (2008).
- 9. B.P. Mandal, S.K. Deshpande and A.K. Tyagi, *J. Mater. Res.*, **23**, 911 (2008).
- A.M. Durand, P. Klavins and L.R. Corruccini, J. Phys. Condens. Matter, 20, 1 (2008).
- 11. B.P. Mandal and A.K. Tyagi, BARC Newssletter, Issue No 313 (2008).
- 12. A.N. Thakur, Ph.D. Thesis, Gorakhpur University, Gorakhpur, India (1992).
- 13. L.F. Bates, Modern Magnetism, Cambridge, London (1951).
- C. Kittel, Introduction to Solid State Physics, Wiley, New York, edn. 7 (1996).
- 15. A.J. Dekker, Solid State Physics, MacMillan, London (1964).
- A.N. Thakur, K. Gaur, M.A. Khan and H.B. Lal, *Indian J. Phys.*, **71A**, 91 (1997).